

Notes for Lectures on Quantum Mechanics

Some Useful Restrictions on CG coefficients

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In this connection with addition of angular momenta, the following results from the theory of angular momentum derived earlier will be useful.

$$J_{\pm}|JM\rangle = \sqrt{(J(J+1) - M(M \pm 1))}|JM \pm 1\rangle \quad (1)$$

$$\begin{aligned} (J_{\pm}^{(1)} + J_{\pm}^{(2)})|j_1 m_1 j_2 m_2\rangle &= \sqrt{(j_1(j_1+1) - (m_1 \pm 1))}|j_1 m_1 \pm 1 j_2 m_2\rangle \\ &+ \sqrt{(j_2(j_2+1) - m_2(m_2 \pm 1))}|j_1 m_1 j_2 m_2 \pm 1\rangle \end{aligned} \quad (2)$$

On taking conjugate of Eq.(2) we get

$$\begin{aligned} \langle j_1 m_1 j_2 m_2 | (J_{\mp}^{(1)} + J_{\mp}^{(2)}) &= \langle j_1(m_1 \pm 1) j_2 m_2 | \sqrt{(j_1(j_1+1) - (m_1 \pm 1))} \\ &+ \langle j_1 m_1 j_2(m_2 \pm 1) | \sqrt{j_2(j_2+1) - m_2(m_2 \pm 1)} \end{aligned} \quad (3)$$

which is a consequence of the angular momentum commutation relations. Considering the matrix element

$$\langle j_1 j_2 m_1 m_2 | J_{\pm} | JM \rangle = \langle j_1 j_2 m_1 m_2 | (J_{\pm}^{(1)} + J_{\pm}^{(2)}) | JM \rangle \quad (4)$$

and using Eq.(1) and Eq.(3) we get two relations, one for J_+ and

$$\begin{aligned} &\sqrt{J(J+1) - M(M+1)} \langle j_1 j_2 m_1 m_2 | J(M+1) \rangle \\ &= \langle j_1 j_2 m_1 - 1 m_2 | JM \rangle \sqrt{j_1(j_1+1) - m_1(m_1+1)} \\ &+ \langle j_1 j_2 m_1 m_2 - 1 | JM \rangle \sqrt{j_2(j_2+1) - m_2(m_2+1)} \end{aligned} \quad (5)$$

and a second relation for J_-

$$\begin{aligned} &\sqrt{J(J+1) - M(M-1)} \langle j_1 m_1, j_2 m_2 | J(M-1) \rangle \\ &= \langle j_1(m_1+1) j_2 m_2 | JM \rangle \sqrt{j_1(j_1+1) - m_1(m_1-1)} \\ &+ \langle j_1 m_1, j_2(m_2+1) | JM \rangle \sqrt{j_2(j_2+1) - m_2(m_2-1)} \end{aligned} \quad (6)$$

We will make repeated use of the results Eq.(5),Eq.(6) given above. These equations can be used successively with $M = J, J-1, \dots$ to compute the Clebsch Gordon coefficients.

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